Binary black-hole spin precession: a tale of three timescales

Davide Gerosa

arXiv:1506.03492

(to be submitted to PRD)

arXiv:1411.0674

(**PRL** 114:081103)

with M.Kesden, U.Sperhake, E.Berti and R.O'Shaughnessy



Department of Applied Mathematics and Theoretical Physics

June 12th, 2015 One Hundred Years of Strong Gravity Lisbon

> d.gerosa@damtp.cam.ac.uk www.damtp.cam.ac.uk/user/dg438

Outline

- 1. A timescale hierarchy
- 2. Analytical precession
- 3. Quasi-adiabatic inspiral
- 4. Morphological phase transitions



Spinning BH binaries?



Three dynamics, and three times

1. Orbital motion

- $t_{
 m orb} \propto (r/r_g)^{3/2}$ Kepler's third law
- 2. Spin & orbital-plane precession $t_{\rm pre} \propto (r/r_g)^{5/2}$
- 3. GW emission and inspiral

Feters & Matthews 1963 if (Post-)Newtonian $r \gg r_g = GM/c^2$: timescale hierarchy



BH binary **multi-timescale** analysis:

- 1. Solve the dynamics (hopefully analytically) on the shorter time
- 2. Quasi-adiabatic evolution ("average") on the longer time

Common practice in binary dynamics

Apostolatos et al 1994

 $t_{\rm RR} \propto (r/r_g)^4$

Quadrupole formula

- periastron precession
- osculating orbital elements
- variation of constants

Usual Post-Newtonian dynamics



Variables

- Total mass (units) $M = m_1 + m_2 = 1$
- Mass ratio $q = m_2/m_1 \le 1$
- Three momenta, 9 components ${f L}, {f S_1}, {f S_2}$

Orbit-average PN evolutionary equations e.g. Kidder 1995

- Spin precession $\dot{\mathbf{S}}_i = \mathbf{\Omega}_i \times \mathbf{S}_i$
- Momentum conservation $\hat{\mathbf{L}} = -(\hat{\mathbf{S_1}} + \hat{\mathbf{S_2}})/\mathbf{L}$
- Radiation reaction $\dot{r} = PN$ approximant $L = q\sqrt{rM^3/(1+q)^2}$

Constraints

- Spin magnitudes are constants $S_i = \chi_i m_i^2$
- Take a smart frame (3 constraints)

Double-spin precession is (usually) a 4D problem: $r, \theta_1, \theta_2, \Delta \Phi$



Inspiral

A closer look...

Orbit

4D is too much. Can we further exploit the **timescale hierarchy**

 $\hat{\mathbf{z}}'$

 \mathbf{L}

 $\Delta \Phi$

 S_2

 $\mathbf{\hat{y}}'$

 $\mathbf{S}_{\perp} = \hat{\mathbf{y}} \times \mathbf{S}$

â

Inspiral

ŷ

<< (Precession)

Let's **freeze** GW emission

- Separation r varies on $t_{
 m RR}$
- Also $J = |\mathbf{L} + \mathbf{S_1} + \mathbf{S_2}|$ vary on t_{RR}
- Effective spin is constant (at least) at 28N!

$$\xi \equiv M^{-2}[(1+q)\mathbf{S}_1 + (1+q^{-1})\mathbf{S}_2] \cdot \hat{\mathbf{L}}_{\theta_1} \theta_2$$

Upshot

- This construction can be done explicitly
- The solution is fully analytical
- Need a smarter frame...
- Chosen parameter is $S = |\hat{\mathbf{S_1}} + \mathbf{S_2}|$

Double-spin precession is (actually) a **1D** problem!

On the shoulders of giants



Kepler's two-body problem What you do:

- One effective particle: 3D
- 3D to 2D problem:
 L is a constant of motion!
- Energy is **constant:** 2D to 1D?
- Effective potential

What you get:

- A lot of understanding
- Solutions are Kepler's orbits
- Phases: bound, unbound

Integrating GMm/r^2 to get a bunch of points along an orbit or... **knowing** that that curve is an ellipse!

Effective potentials for spin precession



What you do:

- Start from 4D problem
- 4D to 2D problem: GW are frozen, r and J are constant,
- Further constant of motion,
 effective spin: 2D to 1D
- Effective potentials for BH binary spin precession

What you get:

- Analytical solutions
- Phases: circulating, librating
- A lot of understanding

Integrating the PN eq. to get a bunch of points on a precession cone or... **knowing** the shape of that cone!



Spin morphologies

How do **Solutions** look like?

Spin tilts $\theta_1, \ \theta_2$

- Kind of boring... monotonic
- Bounded by the effective potentials

Azimuthal projections $\Delta\Phi$

- Three different morphologies
- Boundaries if aligned



Complete classification

DG et al. 2015





Averaging the average

<< (Precession)



m

Let's turn on GW emission

• Quasi-adiabatic approach

Orbit

- Relate solutions at different separations
- Only r (or L) and J vary on $t_{
 m RR}$. Can we relate them?

Usual orbit average

$$\langle X \rangle_{\rm orb} = \frac{\int d\psi \ X \ dt/d\psi}{\int d\psi \ dt/d\psi}$$
 Some parameters for the dynamics (here ψ is Kepler's true anomaly) Orbital period

New precession average

$$\langle X \rangle_{\text{pre}} = \frac{\int dS \ \langle X \rangle_{\text{orb}} \ dt/dS}{\int dS \ dt/dS}$$
 Precession period

A **new** Post-Newtonian approach



How do **Solutions** look like?

- PN evolution is reduced to solving one single ODE
- Computationally, very very **easy**
- Domain can be compactified and integrations carried over from $r/M = \infty$



Morphological phase transitions



 $\hat{}$ At small separations: libration $\,\Delta\Phi\sim 0$ or $\Delta\Phi\sim\pi$

The role of alignment

Evolution of the tilt angles on the inspiral time...

(allowed range of the evolution on the precessional time)

- The **range** grows fatter
- "Bounce" at the alignment configuration...
- ... and sharp transition towards another morphology

 \mathbf{S}

Circulating

 $S_1 S_2$

 $\Delta \Phi$.



A predictive statement

• The morphology **is a feature** of spin precession that does <u>not</u> vary on the precessional time!

Initial means

astrophysics

From the final

morphology,

one estimates

the initial spin

orientation:

how BH form!

- No theoretical uncertainties related to the precessional phase
- The final spin orientations are scattered around, but...
-each blue, green, red dots comes from the blue, green, red region!



A multi-timescale perspective

Usual PN approach:

orbit-averaged equations.

You don't follow each BH along its orbit, but consider the orbits <u>"as a whole"</u>

Our new PN approach:

precession-averaged equations.

You don't even track each spin along its precession cone, but consider the cones <u>"as a whole"</u>



Summary

- 1. A timescale hierarchy
- 2. Analytical precession

5. From **GR** to **astro**

- 3. Quasi-adiabatic inspiral
- 4. Morphological phase transitions

